



# Decay Detector For The Study of Isoscalar Giant Monopole Resonances

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## Giant Resonances

- Collective excitations of the nuclei
- Discovered in the 1940s while bombarding nuclei with gamma rays
- Monopole resonance is a spherical oscillation
- Isoscalar – neutrons and protons move in phase with one another
- Isovector – neutrons and protons move out of phase with one another

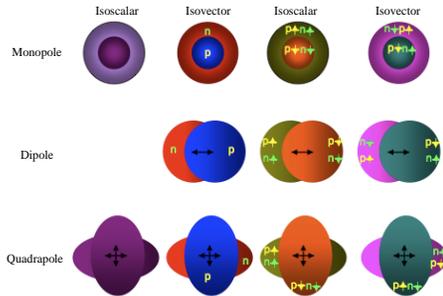


Figure 1: depiction of giant resonance modes (ref. Xinfeng Chen, "Giant Resonance Study By  $^6\text{Li}$  Scattering")

- **Classical description:** Liquid Drop Model
  - Protons and neutrons are treated as independent fluids
  - Giant resonance can be thought of as oscillations of density or shape of these fluids
- **Quantum Mechanical Description:** Nuclear Shell Model
  - Analogous to the electron shell model
  - Neutrons and protons have their own shell system
  - Levels in the system correspond to energy levels that protons and neutrons jump between during various processes
  - Resonance can be thought of by protons and neutrons jumping up to a higher energy level leaving a hole in the lower shell
- Energy Weighted Sum Rule (EWSR) is used to measure the strength of the giant resonance
 
$$S(Q) \equiv \sum_n \langle E_n - E_0 \rangle | \langle n | Q | 0 \rangle |^2 = \frac{1}{2} \langle 0 | [Q, [H, Q]] | 0 \rangle$$

## Purpose

- From studying the Isoscalar Giant Monopole Resonance (ISGMR),  $E_{\text{GMR}}$  can be determined
- $E_{\text{GMR}}$  can be used to find  $K_{nm}$ , a main component in the nuclear matter equation of state
- Astrophysics:
  - Supernova collapse
  - Neutron Stars

## Unstable Nuclei

- Studying ISGMR in unstable nuclei may allow us to determine the dependence of the nuclear incompressibility on  $(N-Z)/A$
- Cannot successfully make a target out of unstable nuclei
- Study the inverse reaction instead
- Normal Kinematics:  $^{26}\text{Si}(^6\text{Li}, ^6\text{Li})^{26}\text{Si}^*$
- Inverse Kinematics:  $^6\text{Li}(^{26}\text{Si}, ^{26}\text{Si}^*)^6\text{Li}$

## Experimental Setup

- Decay particles will hit the decay detector
- Remaining particles will go on to the MDM spectrometer

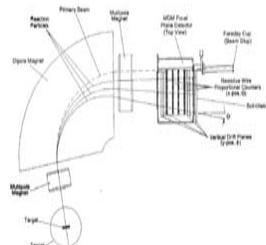


Figure 2: depiction of the experimental setup (ref. Xinfeng Chen, "Giant Resonance Study By  $^6\text{Li}$  Scattering")

- MDM Spectrometer separates particles of different momentum and charge
- Ionization chamber measures the position
- Scintillator measures the total energy

## Decay Detector

- Composed of a layer of horizontal 1 mm thick scintillator strips followed by a layer of vertical 1 mm thick scintillator strips in front of 5 block scintillators

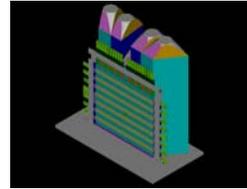


Figure 3: The decay detector (drawing by Dr. Tokimoto)

- Each scintillator is connected to optical fibers that carry the signal to a photomultiplier tube (PMT)
- In the PMT, the photon's energy is converted to an electric signal and is amplified
- Scintillator strips give the angle of the particle as well as the energy loss through the strip
- The strip scintillator is also used for particle ID within the  $7 < E/A < 11$  (MeV/amu)
- The block scintillators measure the remaining energy of the particle
- $11 < E/A < 70$  (MeV/amu) block IDs particles

## Stopping Power

- Stopping power is the energy loss per unit length ( $dE/dx$ )
- Commonly approximated by the Bethe-Bloch formula

$$\frac{dE}{dx} = 2\pi n_0 m_e c^2 n_e Z^2 \left[ \ln \left( \frac{2m_e c^2 \gamma^2 \beta^2 W_{\text{max}}}{I^2} \right) - 2\beta^2 - \delta - 2 \frac{C}{Z} \right]$$

$$\beta = \frac{v}{c}$$

$$\gamma = 1/\sqrt{1-\beta^2}$$

- A more convenient form is

$$\frac{d\varepsilon}{dx} = \frac{Z^2}{m} \frac{1}{(1+\mu)} \frac{\kappa}{(\varepsilon + \varepsilon_0)^\mu}$$

- $\kappa$ ,  $\mu$ , and  $\varepsilon_0$  are medium dependent constants that we fit to SRIM table data
- For our purposes,  $\varepsilon$  is equivalent to  $\beta$

## Predicting Light Output

- Two methods for predicting the light output
- **Birks semi-empirical formula:**
  - Non-linear relationship to the stopping power at low energies

$$\frac{dL}{dx} = \frac{C_0 dE/dx}{1 + C_1 dE/dx}$$

- $C_0$  is the proportion of the molecules that contribute to the light output
- $C_1$  is the proportion of molecules that are quenching sites
- **Energy Deposition By Secondary Electrons (EDSE) Model:**
  - Advantageous for calibrations across numerous ion types
  - Takes into account both ionization density and energy transport concepts
  - Assumes that there exists a "quenching density,"  $\rho_q$ , which defines a "quenching radius,"  $r_q$
  - Below  $r_q$ , the scintillator response is assumed to be constant
  - Above  $r_q$ , the specific detector response,  $dL/dx$  has a linear dependence on the integral of the density of energy deposited by the scattered electrons

$$\rho(r) = \frac{D}{nr^2} \left( 1 - \frac{r}{R_{\text{max}}} \right)^{d+1/n}$$

$$D = N_A \rho_{\text{det}} \left( \frac{Z_{\text{det}}}{A_{\text{det}}} \right) \frac{e^4}{m_e c^2} \frac{Z^*2}{\beta^2}$$

$$\frac{dN_e}{dx} = C \left[ \pi \rho_q(r_q^2) + \int_{r_q}^{R_{\text{max}}} \rho(r) 2\pi r dr \right]$$

$$\frac{dL}{dx} = C \frac{dN_e}{dx} \left( 1 - \mathcal{F} \frac{dN_e/dx}{\mathcal{H} + dN_e/dx} \right)$$

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